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Understanding Raman Scattering in NIF Ignition Experiments

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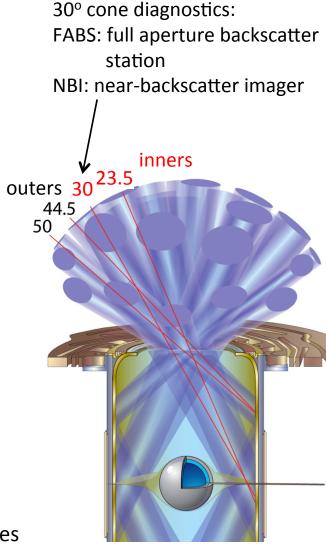
Anomalous Absorption Conference 20 June 1011

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Linear gain analyses better match reflectivity trends with improved plasma and laser beam models

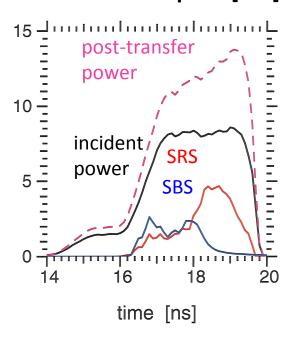
- "High Flux Model" (HFM) for rad-hydro:
 - DCA opacities, 0.15 electron heat flux limiter
 - Cross-beam energy transfer (linear model with clamp)
 - Measured backscatter removed
- Linear gain spectrum with HFM plasma conditions:
 - Close to measured SRS wavelength
 - Agreement better if multi-quad (overlapped beam)
 laser intensity used, rather than single-quad
- Gain and reflectivity time histories:
 - Gain increases in time, while reflectivity first increases and then decreases late in peak power
- Spatially non-uniform cross-beam energy transfer:
 - Gain decreases late in peak power, like measured reflectivity
- **Electron trapping:** pF3D simulations give SRS Langmuir waves above threshold for trapping nonlinearities



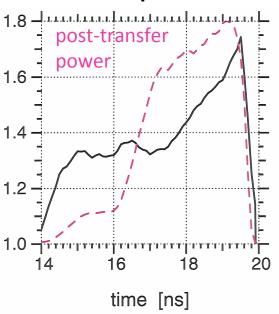
We study NIF shot N110214 - symmetry capsule (symcap) with ~1.3 MJ laser energy - 30° (inner) cone

- "Post-transfer" reflectivity = measured SRS / Lasnex power w/ cross-beam transfer.
- SRS energy reflectivity [joules out / joules in]:
 - Incident power: 27%
 - Post-transfer power: 19%

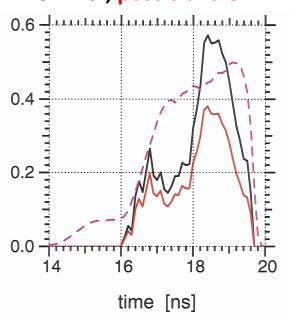
Power on one quad [TW]



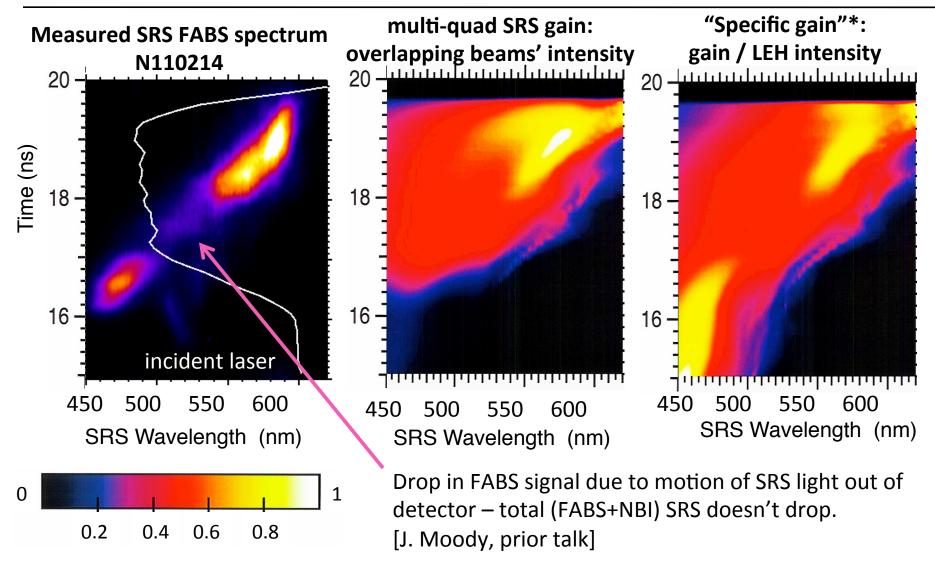
post-transfer power incident power



SRS reflectivity: nominal, post-transfer



SRS wavelength increases in time, indicating SRS occurs at progressively higher density

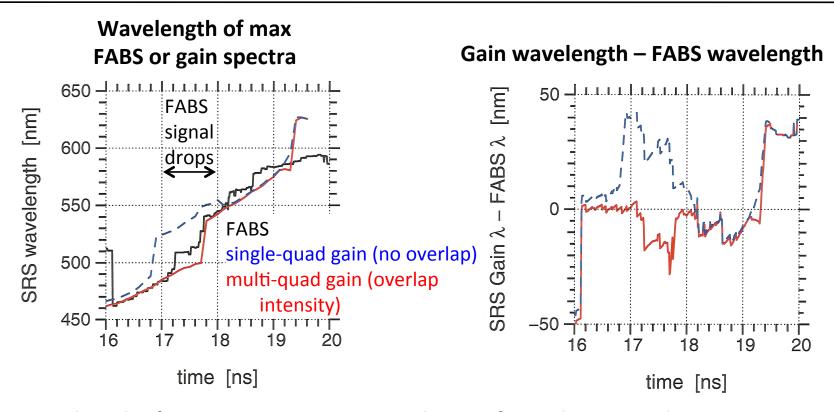


FABS = full aperture backscatter station NBI = near-backscatter imager

*Suggested by L. Suter

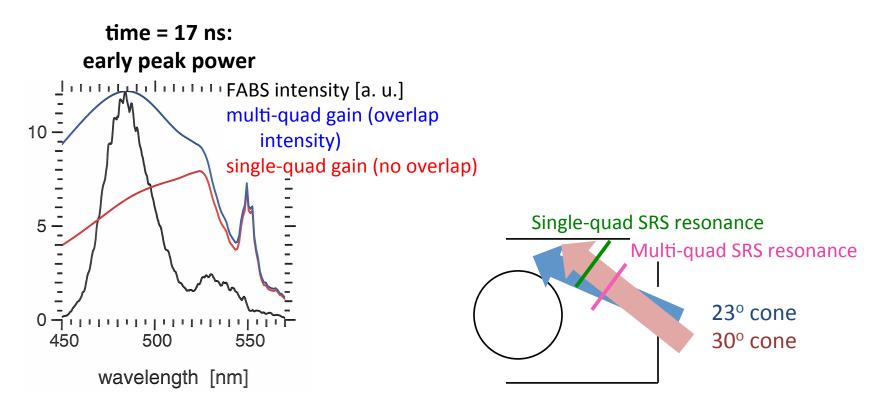
DJS: AA SRS 2011; p. 4

With HFM (high-flux model), wavelength of peak multi-quad gain agrees well with FABS, indicating plasma conditions are ~ right



- Wavelength of max SRS separates power history from plasma conditions.
- Early peak power: 17-18ns:
 - Single-quad gains peak at a longer wavelength.
 - FABS signal reduced due to motion of SRS.
 - longer wavelengths refract more, so FABS light may be shorter wavelength than total SRS light.

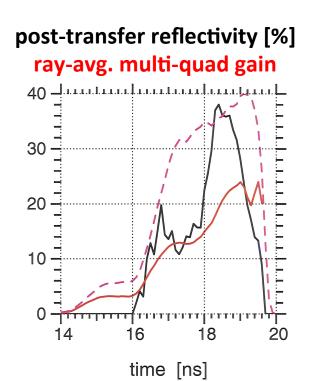
Multi-quad SRS gains peak at shorter wavelength since beams overlap near laser entrance hole

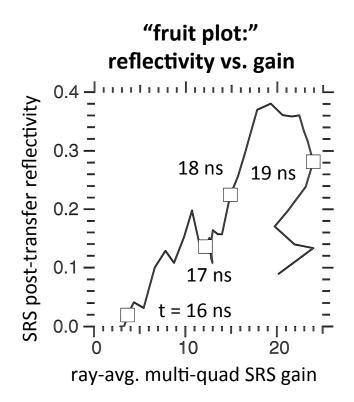


Multi-quad SRS gains peak at shorter wavelength than single-quad gains:

beams overlap more near the LEH, where the electron density, and plasma frequency, is lower.

Spatially uniform transfer: reflectivity scales with gain until late in peak power



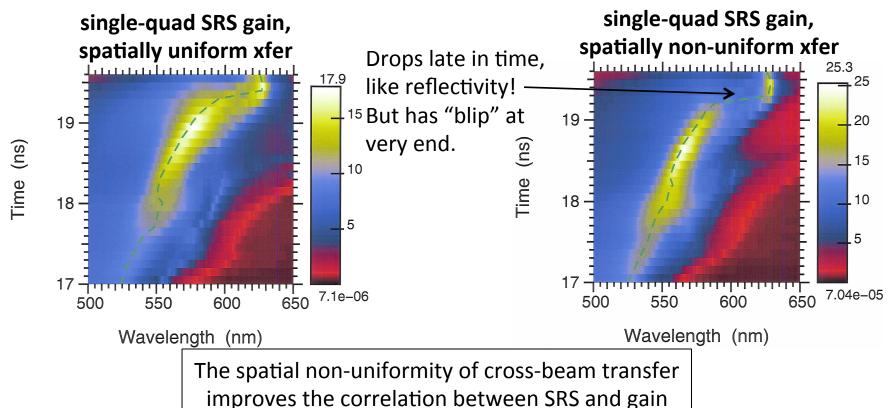


---- post-transfer power

- Gain tracks reflectivity until ~ 18.5 ns (mid-late peak power).
- At late time, reflectivity drops but gain doesn't.
- Late-time gain coming from long wavelengths generally not observed in FABS.

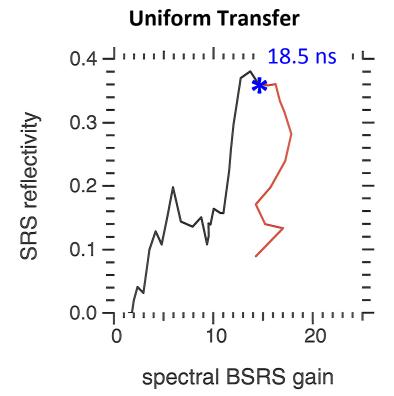
More detailed calculations of cross-beam transfer introduce spatial non-uniformity in the intensities

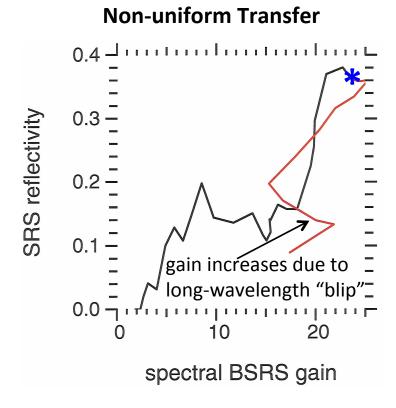
- Current HFM: distributes transferred power uniformly across the beam
- Account for spatial non-uniformity: run SLIP at one time (18 ns): E. A. Williams, later talk
- Provides 3D spatial beam intensity multiplier. Use this mask at all times
- Calculate SRS gain with spatially non-uniform transfer, and single intensities



DJS: AA SRS 2011; p. 8

SRS gain with spatially non-uniform beam transfer tracks reflectivity better than with uniform transfer

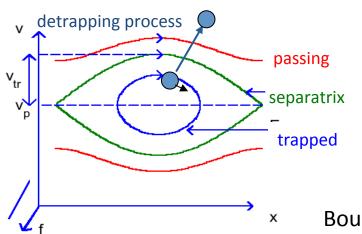




Threshold for electron trapping nonlinearity given by "bounce number" N_R

Electron trapping nonlinearities:

- Landau damping reduction, frequency shift, Langmuir-wave self-focusing.
- Effective only if electrons resonant w/ Langmuir wave complete ~ 1 bounce orbit before being detrapped.



Important detrapping processes:

- 1. Speckle sideloss (geometric effect): N_{B,sl}
- Collisions: electron-electron and electron-ion treated together: N_{B,coll}

(SSD way too slow to matter)

Bounce number:
$$N_B \equiv \frac{\tau_{de}}{\tau_B} = \frac{\text{detrapping time}}{\text{bounce period}} \sim \delta n^{1/2}$$

sideloss:
$$N_{B,sl} = \left[\frac{\delta n}{\delta n_{sl}}\right]^{1/2}$$
 $\delta n_{sl} = 2.67 \left(\frac{8}{F} \frac{\lambda_{De}}{\lambda_0}\right)^2$ Bounce period: $\tau_B = \frac{2\pi}{\omega_{ne}} \sqrt{\frac{n_e}{\delta n}}$

$$\delta n_{\rm sl} = 2.67 \left(\frac{8}{F} \frac{\lambda_{De}}{\lambda_0} \right)^2$$

$$\tau_{\scriptscriptstyle B} \equiv \frac{2\pi}{\omega_{\scriptscriptstyle pe}} \sqrt{\frac{n_{\scriptscriptstyle e}}{\delta n}}$$

collisions: $N_{B,coll} = \left[\frac{\delta n}{\delta n_{coll}} \right]^{3/2}$

$$N_{R}^{-1} = N_{R,cl}^{-1} + N_{R,coll}^{-1}$$

Joint bounce number: $N_B^{-1} = N_{B,sl}^{-1} + N_{B,coll}^{-1}$ Independent detrapping rates add.

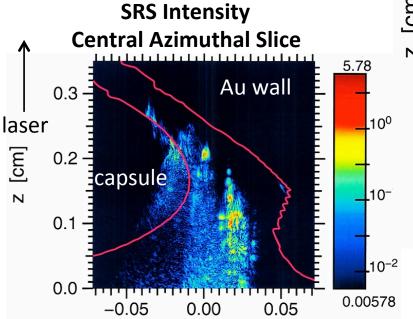
Trapping assessment of pF3D run suggests trapping occurs in parts of the 30 degree beam

pF3D:

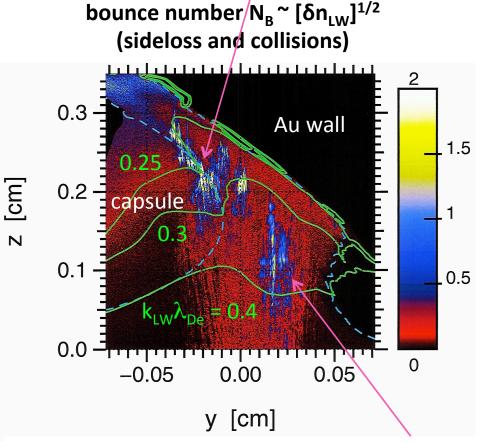
- parallel, paraxial envelope code
- linear plasma response used

N110214 profiles, time = 18 ns:

- Trapping can change the local gain
- pF3D SRS reflectivity ~ 20%



Risk of trapping or Langmuir Decay Instability



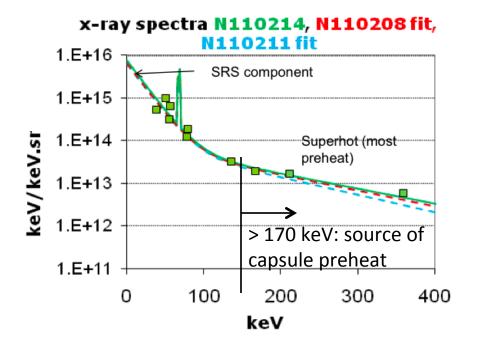
Conclusions

- "High Flux Model" rad-hydro with spatially uniform cross-beam energy transfer:
 - SRS gain spectrum agrees well with measurements
 - Especially when multi-quad laser intensity used
 - Except for long-wavelength gains late in time not seen in experiments
- Spatially uniform transfer: reflectivity and gain correlate until late in peak power
 - Late in time, reflectivity drops but gain does not
- Spatially non-uniform transfer: the correlation of reflectivity and gain improves
- Electron trapping: pF3D simulation shows regions in the beam where Langmuir wave amplitudes above threshold
 - May play a role in some of the SRS seen

A cross-beam transfer model, including spatial non-uniformity and plasma profile modification, is being added to Hydra

Backup slides after here

Observed SRS is consistent with the colder component of the hot electron spectrum



$$\frac{dN}{dE} \sim \frac{E_1}{T_1^2} \exp[-E/T_1] + \frac{E_2}{T_2^2} \exp[-E/T_2]$$

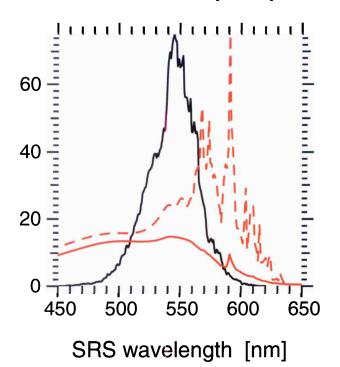
"SRS component":

$$E_1$$
 = 70 kJ, T_1 = 18 keV
0.5 m_e v²_{ph,LW} = 18 keV for λ_{SRS} = 570 nm
"conventional" backward SRS, measured in FABS/NBI

"Superhot component": $E_1 = 0.8 \text{ kJ}$, $T_1 = 124 \text{ keV}$ independent LPI process, such as: two-plasmon decay, backward SRS at $n_{crit}/4$, forward SRS, ...

Significant gain can occur at longer wavelengths than measured on FABS

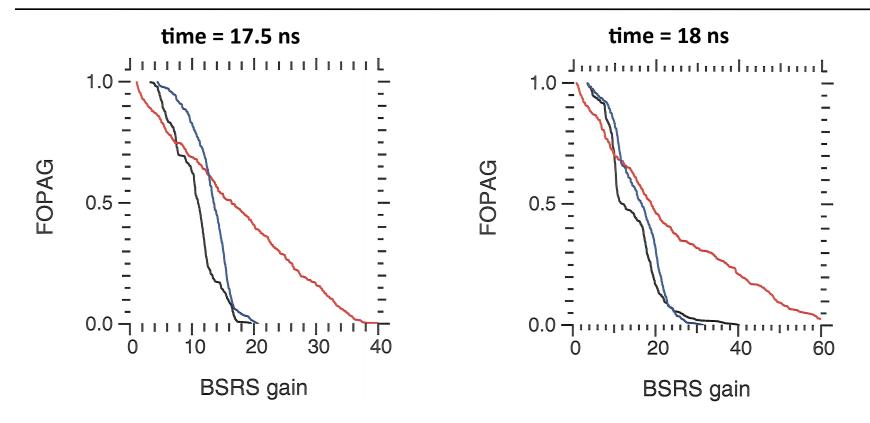
time = 18 ns – mid peak power



FABS intensity [a. u.]
——— multi-quad ray-avg. gain
---- multi-quad log [ray-avg. exp (G)]

- Long-wavelength SRS washed out in rayaveraging, since each ray has a narrow peak (weak damping) at a different wavelength.
- Long-wavelength SRS may not occur: shorterwavelength SRS occurs at lower density, nearer the LEH, and may deplete the pump.
- If it does occur, it will be more refracted than shorter-wavelength light [c.f. J. Moody's talk] and may miss the FABS detectors.
- Also, it will be more absorbed in the target by inverse bremsstrahlung.

Spatially varying beam transfer gives a wider distribution of ray gains

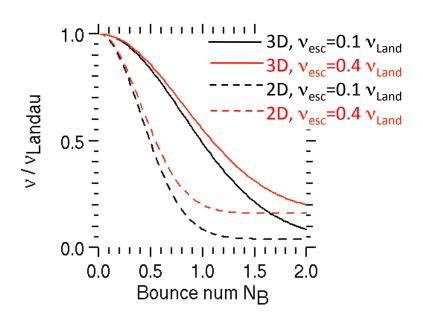


single-quad gain, uniform xfer multi-quad gain, uniform xfer single-quad gain, varying xfer FOPAG = fraction of ray power above a gain.

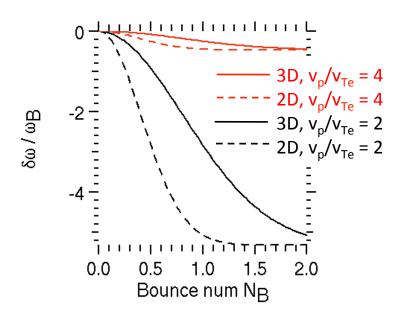
For each ray: find the max gain within λ = +- 10 nm of $<\lambda_{max}>_{avg}$.

Damping reduction and frequency shift in finite-radius Langmuir wave: theory by H. A. Rose

Damping Reduction: more rapid in 2D than 3D



Frequency Downshift: rapidly increases with $k\lambda_D$



Fraction of coupling above a bounce number: allows quantification of trapping

$$\frac{dE_{scat}}{dt} \mid_{coup} \propto E_{las} \delta n_{epw}$$

63.9 ps [black], 65.5 ps [red], sideloss + collisions [solid], sideloss [dash]

